

THE TIME TRANSFER AND SYNGE WORLD FUNCTIONS IN RPS

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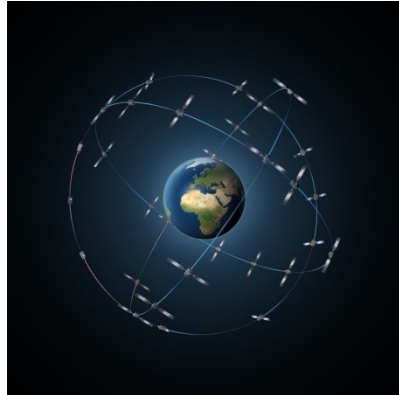
Outline

- Motivation
- Timing corrections
- Time transfer function
- Synge world function
- Simulating null coordinates
- Solving inverse problem
- Still to come

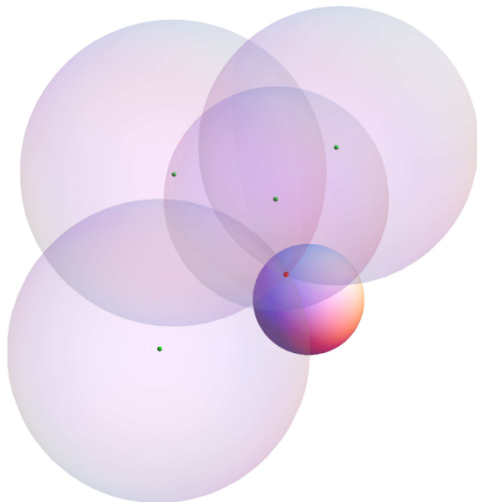
Motivation

GNSS – Global navigation satellite system

GPS,
GLONASS,
GALILEO,
BEIDOU



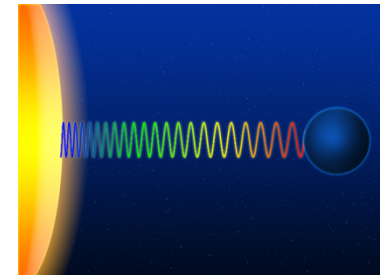
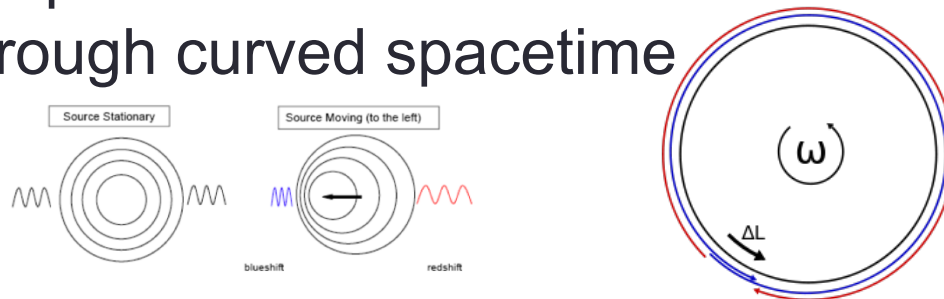
$$c^2(t_P - t_i)^2 = (x_i - x_P)^2 + (y_i - y_P)^2 + (z_i - z_P)^2$$



- RPS – a fully relativistic positioning system
- High accuracy with no corrections
- Autonomy – no ground tracking required
- Enables measurements of gravitational field

Timing corrections

- Proper time(τ): time indicated by an on-board clock
- Coordinate time (t): time indicated by a clock very far away from any mass
- Requirement: to determine the time of flight of light rays through curved spacetime



- Time transfer and Synge world function relationship

$$\sigma(t, x^i, t', x^{i'}) = \frac{1}{2}(\lambda_1 - \lambda_0) \int_{\lambda_0}^{\lambda_1} g_{ab}(x) \frac{dx^a}{d\lambda} \frac{dx^b}{d\lambda} d\lambda$$

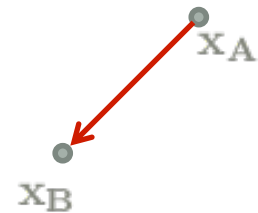
$$\sigma(t, x^i, t - \mathcal{T}(x^i, x^{i'}), x^{i'}) = 0$$

Time transfer function

- Time transfer function

$$\frac{1}{c} \int_{x_B}^{x_A} dx^0 = t_B - t_A = \mathcal{T}_e(t_A, \mathbf{x}_A, \mathbf{x}_B) = \mathcal{T}_r(\mathbf{x}_A, t_B, \mathbf{x}_B)$$

$$\mathcal{T}(\mathbf{x}_A, \mathbf{x}_B) = \frac{|\mathbf{x}_A - \mathbf{x}_B|}{c} + \sum_{n=1}^{\infty} \mathcal{T}^{(n)}(\mathbf{x}_A, \mathbf{x}_B)$$



- Post Minkowskian metric

$$ds^2 = \mathcal{A}(r)(dx^0)^2 - \mathcal{B}^{-1}(dr^2 + r^2 d\theta^2 + r^2 \cos^2 \theta d\phi^2)$$

$$\mathcal{A}(r) = 1 - 2\frac{m}{r} + 2\beta\frac{m^2}{r^2} - \frac{3}{2}\beta_3\frac{m^3}{r^3} + \beta^4\frac{m^4}{r^4} + \sum_{n=5}^{\infty} \frac{(-1)^n n}{2^{n-2}} \beta^n \frac{m^n}{r^n}$$

$$\mathcal{B}^{-1}(r) = 1 + 2\gamma\frac{m}{r} + \frac{3}{2}\epsilon\frac{m^2}{r^2} + \frac{1}{2}\gamma_3\frac{m^3}{r^3} + \frac{1}{16}\gamma_4\frac{m^4}{r^4} + \sum_{n=5}^{\infty} (\gamma_n - 1) \frac{m^n}{r^n}$$

$$U(r) = \frac{1}{\mathcal{A}(r)\mathcal{B}(r)} = 1 + 2(1 + \gamma)\frac{m}{r} + \sum_{n=2}^{\infty} 2\kappa_n \frac{m^n}{r^n}$$

- Quasi Minkowskian light rays

$$\begin{cases} x^0 = ct_a + \xi|\mathbf{x}_B - \mathbf{x}_A| + \sum_{n=1}^{\infty} X_{(n)}^0(\mathbf{x}_A, \mathbf{x}_B, \xi) \\ \mathbf{x} = \mathbf{z}(\xi) + \sum_{n=1}^{\infty} \mathbf{X}_{(n)}(\mathbf{x}_A, \mathbf{x}_B, \xi) \end{cases}$$

Time transfer function

- Solving for dt

$$cdt = \frac{J}{E}d\phi \pm \frac{1}{r}\sqrt{r^2\mathcal{U}(r) - \left(\frac{J}{E}\right)^2} dr$$

- Time transfer function

$$c\mathcal{T}(r_A, r_B, \mu) = \frac{J}{E}(\phi_A - \phi_B) + \int_{r_A}^{r_B} \frac{1}{r}\sqrt{r^2\mathcal{U}(r) - \left(\frac{J}{E}\right)^2} dr$$

Linet &
Teyssandier, 2014

- Expanded terms

$$\frac{J}{E} = r_c \left[1 + \sum_{n=1}^{\infty} \left(\frac{m}{r_c}\right) q_n \right] \quad r_c = \frac{r_A r_B}{R_{AB}} \sin(\phi_A - \phi_B)$$

$$q_n = -c \left(\frac{r_c}{m}\right)^n \frac{\sin(\phi_A - \phi_B)}{r_c} \frac{\partial \mathcal{T}^{(n)}(r_A, r_B, \mu)}{\partial(\cos \phi_A - \phi_B)}$$

$$R_{AB} = \sqrt{r_A^2 + r_B^2 - 2r_A r_B \cos(\phi_A - \phi_B)}$$

Time transfer function

- The time transfer function

$$\mathcal{T}(\mathbf{x}_A, \mathbf{x}_B) = \frac{|\mathbf{x}_A - \mathbf{x}_B|}{c} + \sum_{n=1}^{\infty} \mathcal{T}^{(n)}(\mathbf{x}_A, \mathbf{x}_B)$$

$$\mathcal{T}^{(1)}(r_A, r_B, (\phi_A - \phi_B)) = m \frac{1 + \gamma}{c} \ln \left(\frac{r_A + r_B + R_{AB}}{r_A + r_B - R_{AB}} \right) \longrightarrow \text{Shapiro, 1964}$$

$$\mathcal{T}^{(2)}(r_A, r_B, (\phi_A - \phi_B)) = \frac{m^2 R_{AB}}{c r_A r_B} \left[\kappa_2 \frac{(\phi_A - \phi_B)}{\sin(\phi_A - \phi_B)} - \frac{(1 + \gamma)^2}{1 + \cos(\phi_A - \phi_B)} \right], \longrightarrow \text{Linet \& Teysandier, 2014}$$

$$\mathcal{T}^{(3)}(r_A, r_B, (\phi_A - \phi_B)) = \frac{m^3}{c} \frac{R_{AB}(r_A + r_B)}{r_A^2 r_B^2 (1 + \cos(\phi_A - \phi_B))} \left[\kappa_3 - (1 + \gamma) \kappa_2 \frac{(\phi_A - \phi_B)}{\sin(\phi_A - \phi_B)} + \frac{(1 + \gamma)^3}{1 + \cos(\phi_A - \phi_B)} \right]$$

$$\begin{aligned} \mathcal{T}^{(4)}(r_A, r_B, \mu) = & \frac{m^4}{c} \frac{R_{AB}}{r_A^3 r_B^3 (1 - \mu^2)^2} \left[-(1 + \gamma)^4 \frac{5(2r_A^2 + 2r_B^2 + r_A r_B(3 - \mu))(1 - \mu)^2}{6(1 + \mu)} \right. \\ & + \kappa_2 (1 + \gamma)^2 \frac{(2r_A^2 + 2r_B^2 + r_A r_B(3 - \mu))(1 - \mu)^2 \arccos \mu}{\sqrt{1 - \mu^2}} \\ & - \kappa_2^2 \frac{\arccos \mu \left(R_{AB}^2 (1 - \mu^2) - (r_B - r_A \mu)(r_B \mu - r_A) \sqrt{1 - \mu^2} \arccos \mu \right)}{2\sqrt{1 - \mu^2}} \\ & + (1 + \gamma) \kappa_3 \left((1 - \mu) \left((r_A^2 + r_B^2) (3 - \mu) + 2r_A r_B (1 - 3\mu) \right) + R_{AB}^2 \sqrt{1 - \mu^2} \arccos \mu \right) \\ & \left. + \frac{\kappa_4}{2} \left((2r_A r_B - (r_A^2 + r_B^2) \mu) (1 - \mu^2) + R_{AB}^2 \sqrt{1 - \mu^2} \arccos \mu \right) \right] \end{aligned}$$

Synge world function

$$\sigma(x, x') = \frac{1}{2}(\lambda_1 - \lambda_0) \int_{\lambda_0}^{\lambda_1} g_{ab}(z) \frac{dz^a}{d\lambda} \frac{dz^b}{d\lambda} d\lambda$$

$$= \frac{1}{2} \xi (\lambda_1 - \lambda_0)^2$$

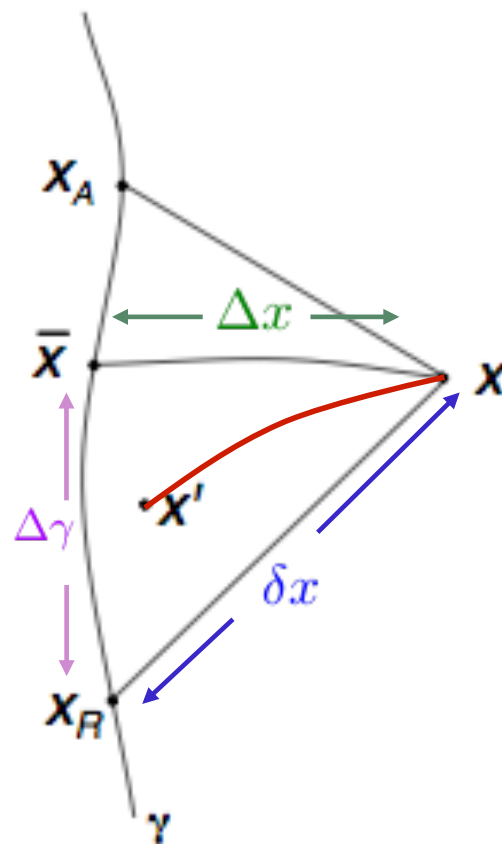
\Rightarrow **for light** $\sigma = 0$

$$g^{ab} \sigma_a \sigma_b = 2\sigma \quad g^{a'b'} \sigma'_a \sigma'_b = 2\sigma$$

$$\sigma = g_{a'b'} \delta x^{a'} \delta x^{b'} \epsilon^2 + A_{a'b'c'} \delta x^{a'} \delta x^{b'} \delta x^{c'} \epsilon^3 + B_{a'b'c'd'} \delta x^{a'} \delta x^{b'} \delta x^{c'} \delta x^{d'} \epsilon^4 + O(\epsilon^5)$$

$$\gamma(\tau') = \gamma(\bar{\tau}) + u(\bar{\tau})(\tau' - \bar{\tau})\epsilon + \dot{u}(\bar{\tau}) \frac{\Delta\tau^2}{2!} \epsilon^2 + O(\epsilon^3)$$

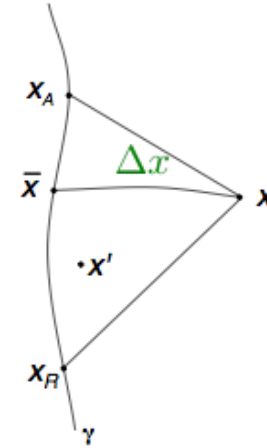
$$\Delta\tau = (\tau' - \bar{\tau}) = \tau_1 + \tau_2 \epsilon + \tau_3 \epsilon^2 + O(\epsilon^3)$$



Syngge world function

- Proper time

$$\tau' = \bar{\tau} + \tau_1 + \tau_2 + \dots$$



$$\tau_1 = \frac{\bar{r} \Delta r u^r}{2m - \bar{r}} + \Delta t E - \Delta \phi L + \rho$$

$$\begin{aligned} \tau_2 = & -\frac{\Delta r^3 m (E^2 \bar{r}^3 - L^2 \bar{r} + 2L^2 m)}{2\rho \bar{r}^2 (\bar{r} - 2m)^3} + \Delta r \left(\frac{\Delta \theta^2 (E^2 \bar{r}^3 - L^2 \bar{r} + 4m \bar{r}^2 - 2\bar{r}^3 + 2L^2 m)}{4m\rho \bar{r} - 2\rho \bar{r}^2} \right. \\ & \left. + \frac{\Delta \phi^2 (E^2 \bar{r}^3 - 3L^2 \bar{r} + 4m \bar{r}^2 - 2\bar{r}^3 + 6L^2 m)}{4m\rho \bar{r} - 2\rho \bar{r}^2} - \frac{\Delta \phi L}{\bar{r}} \right) \\ & + \Delta t^2 \left(\frac{\Delta r m ((2 - 3E^2) \bar{r}^3 + L^2 \bar{r} - 4m \bar{r}^2 - 2L^2 m)}{2\rho \bar{r}^4 (2m - \bar{r})} + \frac{\Delta \phi L m u^r}{2\rho \bar{r}^2} - \frac{m u^r}{2\bar{r}^2} \right) \\ & + \Delta t \left(\Delta r \left(\frac{\Delta \phi E L (\bar{r} - m)}{\rho \bar{r} (2m - \bar{r})} + \frac{E m}{r (r - 2m)} \right) \frac{\Delta r^2 E m u^r}{2\rho (\bar{r} - 2m)^2} + \frac{\Delta \theta^2 E \bar{r} u^r}{2\rho} + \frac{\Delta \phi^2 E \bar{r} u^r}{2\rho} \right) - \frac{\Delta t^3 E m u^r}{2\rho \bar{r}^2} \\ & + \Delta r^2 \left(\frac{\Delta \phi L u^r (2\bar{r} - 5m)}{2\rho (\bar{r} - 2m)^2} + \frac{m u^r}{2(\bar{r} - 2m)^2} \right) + \Delta \theta^2 \left(\frac{\bar{r} u^r}{2} - \frac{\Delta \phi L \bar{r} u^r}{2\rho} \right) - \frac{\Delta \phi^3 L \bar{r} u^r}{2\rho} + \frac{1}{2} \Delta \phi^2 \bar{r} u^r \end{aligned}$$

$$\rho^2 = (g_{ab} \Delta x^a u^b)^2 + g_{ab} \Delta x^a \Delta x^b$$

Simulating null coordinates

- Post Minkowskian to Schwarzschild

$$\begin{aligned}
 t' &= t, \\
 r' &= r \left[1 - \frac{r_S}{2r} - \sum_{n=2}^{\infty} \left(\frac{r_S}{2^n r} \right)^n \right], \\
 \theta' &= \theta, \\
 \phi' &= \phi.
 \end{aligned}$$

- Baby test

$$\begin{aligned}
 r_0 &= 4.2 \times 10^7 \text{ m}, & r_P &= 5.0 \times 10^7 \text{ m}, \\
 \phi_0 &= 0, & \theta_P &= \pi/2, \\
 t_0 &= 0. & \phi_P &= 0.
 \end{aligned}$$

- Results

Order	$t_p = 1s \Rightarrow \tau$ [s]
Delva	0.9733148699
TT(1)	0.9733148699934330221...
TT(3)	0.9733148699934331269...
TT(4)	0.9733148699934331269...
SW(2)	0.9733148699934928242...
SW(5)	0.9733148699934331214...
SW(6)	0.9733148699934331279...

- Schwarzschild circular orbit

$$\begin{aligned}
 t(\tau) &= \frac{c}{\sqrt{1 - \frac{3r_S}{r_0}}} \tau + t_0, \\
 r(\tau) &= r_0, \\
 \theta(\tau) &= \frac{\pi}{2}, \\
 \phi(\tau) &= \sqrt{\frac{c^2 r_S}{2r_0^3 \left(1 - \frac{3r_S}{2r_0}\right)}} \tau + \phi_0
 \end{aligned}$$

Delva & Olympio 2011

Order	$t_p = 1000s \Rightarrow \tau$ [s]
Delva	999.9710561425
TT(3)	999.97105615857394...
TT(4)	999.971056183339852...
SW(5)	999.971056071208982...
SW(6)	999.971056071208977...

Simulating null coordinates

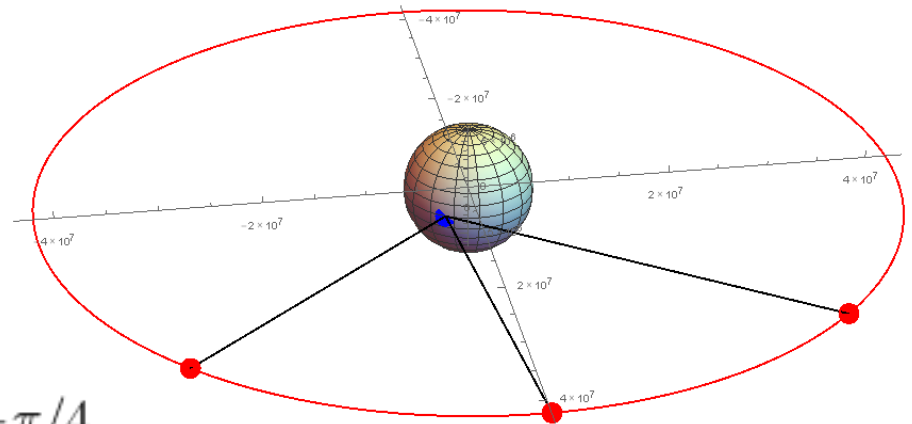
- Three satellites

$$x_P = (1, 6.3 \times 10^6, \pi/2, -\pi/6)$$

$$t_{0,i} = 0 \quad \forall i,$$

$$r_{0,i} = 4.2 \times 10^7 \quad \forall i,$$

$$\phi_{0,1} = 0, \quad \phi_{0,2} = \pi/4, \quad \phi_{0,3} = -\pi/4$$



- Results

Order	Satellite 1: $t_P = 1s \Rightarrow \tau$ [s]
TT(4)	0.8776494176161300...
SW(12)	0.8776494176142050...
SW(11)	0.8776494175963577...

Order	Satellite 2: $t_P = 1s \Rightarrow \tau$ [s]
TT(4)	0.86381940526182...
SW(12)	0.86381940572704...
SW(11)	0.86381940622337...

Order	Satellite 3: $t_P = 1s \Rightarrow \tau$ [s]
TT(4)	0.8800785714457471...
SW(12)	0.8800785714471294...
SW(11)	0.8800785714474831...

Solving inverse problem

- Expand everything!

$$\begin{aligned}
 t_{P,exp} &= t_P^{(0)} + \sqrt{\epsilon} t_P^{(1/2)} + \epsilon t_P^{(1)} + \dots + \epsilon^{order} t_P^{(order)}, \\
 r_{P,exp} &= r_P^{(0)} + \sqrt{\epsilon} r_P^{(1/2)} + \epsilon \left(r_P^{(1)} - m \right) + \mathcal{O}(\epsilon^{3/2}) \\
 \phi_{P,exp} &= \phi_P^{(0)} + \sqrt{\epsilon} \phi_P^{(1/2)} + \epsilon \phi_P^{(1)} + \dots + \epsilon^{order} \phi_P^{(order)}.
 \end{aligned}$$

$$t_{i,exp} = \tau + \frac{3m}{2r_0} \tau \epsilon + \mathcal{O}(\epsilon^2),$$

$$r_{i,exp} = r_0 - m\epsilon + \mathcal{O}(\epsilon^2),$$

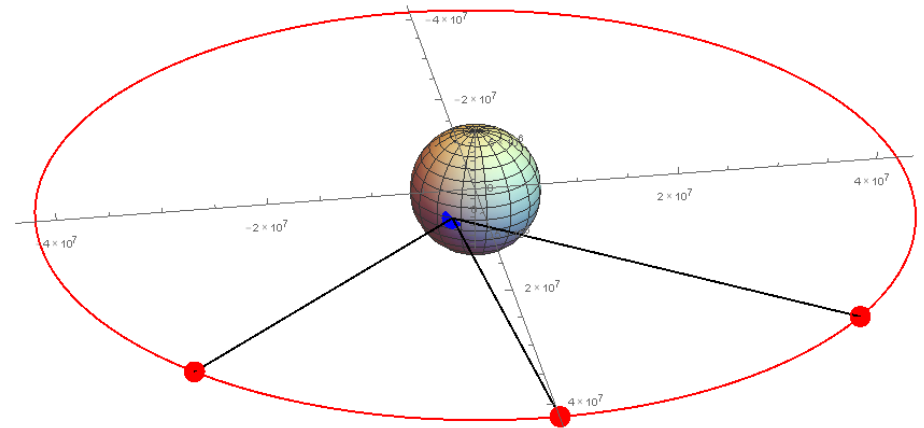
$$\phi_{i,exp} = \phi_{i,0} + \sqrt{\frac{m}{r_0^3}} \tau \sqrt{\epsilon} + \sqrt{\frac{m^3}{4r_0^5}} 3\tau \epsilon^{3/2} + \mathcal{O}(\epsilon^{5/2})$$

$$\begin{aligned}
 \mu_{i,exp} &= \cos \left(\phi_i - \phi_P^{(0)} \right) + \sin \left(\phi_{i,0} - \phi_P^{(0)} \right) \left(\phi_P^{(1/2)} - \sqrt{\frac{m}{r_0^3}} c\tau_i \right) \sqrt{\epsilon} \\
 &\quad - \left[\cos \left(\phi_{i,0} - \phi_P^{(0)} \right) \frac{\left(\sqrt{\frac{4m}{r_0}} c\tau_i - 2r_0 \phi_P^{(1/2)} \right)^2}{8r_0^2} - \sin \left(\phi_{i,0} - \phi_P^{(0)} \right) \phi_P^{(1)} \right] \epsilon \\
 &\quad + \left[\frac{1}{6} \sin \left(\phi_{i,0} - \phi_P^{(0)} \right) \left(\sqrt{\frac{m}{r_0^3}} c\tau_i - \phi_P^{(1/2)} \right)^3 + \cos \left(\phi_{i,0} - \phi_P^{(0)} \right) \left(\sqrt{\frac{m}{r_0^3}} c\tau_i - \phi_P^{(1/2)} \right) \phi_P^{(1)} \right. \\
 &\quad \left. + \sin \left(\phi_{i,0} - \phi_P^{(0)} \right) \left(\phi_P^{(3/2)} - \frac{3 \left(\frac{2m}{r_0} \right)^{3/2} c\tau_i}{4\sqrt{2}r_0} \right) \right] \epsilon^{3/2} + \mathcal{O}(\epsilon^2)
 \end{aligned}$$

Solving the inverse problem

- 3 Satellites

Order	t_P [s]	Error
Actual	1.0	
0	1.000000031160543317958931839335	3.1161×10^{-8}
0.5	0.99999999801722208034042239405	1.9828×10^{-10}
1	1.000000000000000066805491440617	6.6805×10^{-17}
1.5	1.0000000000000000228863072788	2.2886×10^{-19}
2	0.999999999999999999999718036	2.8196×10^{-25}
2.5	1.0000000000000000000000133	1.3263×10^{-28}
3	1.000000000000000000000000000	9.4083×10^{-34}

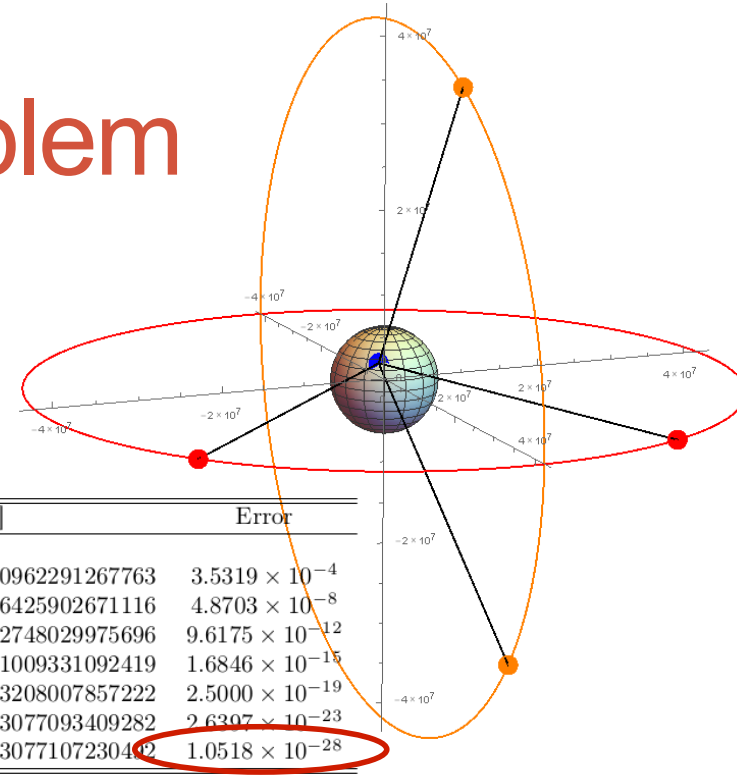


Order	r_P [10^6 m]	Error
Actual	6.3	
0	6.299989882179941576678355635747	-1.6060×10^{-6}
0.5	6.30000000935775873822418454687	1.4854×10^{-10}
1	6.29999999999993758182253996945	9.9076×10^{-16}
1.5	6.2999999999999943035047465108	9.0421×10^{-18}
2	6.30000000000000000000025068941	3.9792×10^{-24}
2.5	6.29999999999999999999973536	4.2006×10^{-27}
3	6.300000000000000000000000000	4.7364×10^{-32}

Order	ϕ_P [rad]	Error
Actual	$-\pi/6$	
0	-0.523663456801297482319083172952	1.1235×10^{-4}
0.5	-0.523598775560885789129469487053	-7.1454×10^{-11}
1	-0.523598775598309328734082371606	1.9969×10^{-14}
1.5	-0.523598775598298889248115917379	3.0884×10^{-17}
2	-0.523598775598298873077052496507	1.0453×10^{-22}
2.5	-0.523598775598298873077107223789	1.2906×10^{-26}
3	-0.523598775598298873077107230547	1.6206×10^{-32}

Solving the inverse problem

- Inclined orbits – 4 satellites



Order	θ_P [rad]	Error
Actual	$\pi/3$	
0	1.047149706632208449047833646793	04.5688×10^{-5}
0.5	1.047197570905738944132225900057	1.8821×10^{-8}
1	1.047197551194164173407845979193	2.3239×10^{-12}
1.5	1.047197551196597732537859875917	1.3003×10^{-17}
2	1.047197551196597746256002605638	9.7201×10^{-20}
2.5	1.047197551196597746154180690784	3.2248×10^{-23}
3	1.047197551196597746154214468106	6.6967×10^{-27}

Order	ϕ_P [rad]	Error
Actual	$-\pi/6$	
0	-0.523783708168425960962291267763	3.5319×10^{-4}
0.5	-0.523598750097320206425902671116	4.8703×10^{-8}
1	-0.523598775603334562748029975696	9.6175×10^{-12}
1.5	-0.523598775598297991009331092419	1.6846×10^{-15}
2	-0.523598775598298873208007857222	2.5000×10^{-19}
2.5	-0.523598775598298873077093409282	2.6307×10^{-23}
3	-0.523598775598298873077107230492	1.0518×10^{-28}

Order	r_P [10^6 m]	Error
Actual	6.3	
0	6.297991095761440172150549624619	3.1887×10^{-4}
0.5	6.299999628817188546413163437225	5.8918×10^{-8}
1	6.29999999954395965929765837974	7.2387×10^{-12}
1.5	6.29999999999990124720349620199	1.5675×10^{-15}
2	6.299999999999998562785122411	2.2813×10^{-19}
2.5	6.2999999999999999734428831	4.2154×10^{-23}
3	6.29999999999999999999941663	9.2599×10^{-27}

Order	t_P [s]	Error
Actual	1.0	
0	1.000004726525898470419881664511	4.7265×10^{-6}
0.5	1.000000000317128497477299209190	3.1713×10^{-10}
1	1.00000000000068448219616118322	6.8448×10^{-14}
1.5	1.00000000000000004276567159170	4.2766×10^{-18}
2	1.00000000000000000000289690092	2.8969×10^{-22}
2.5	0.999999999999999999999890100	1.0990×10^{-25}
3	0.9999999999999999999999967	3.3013×10^{-29}

Solving the inverse problem

- Elliptic orbits – equations of motion

$$\tau = \frac{Y}{2\pi} (1 - e^2)^{3/2} \sqrt{1 - \frac{m}{l}(3 + e^2)} \int d\psi \left\{ \left(\frac{1 - e^2}{1 - e \cos \psi} \right)^{-2} \left(1 - 2 \frac{m}{l} \left(2 + \frac{1 - e^2}{1 - e \cos \psi} \right) \right)^{-1/2} \frac{\sqrt{1 - e^2}}{1 - e \cos \psi} \right\}$$

$$r(\psi) = \frac{l}{1 + e \frac{e - \cos \psi}{e \cos \psi - 1}}$$

$$t = \frac{Y}{2\pi} (1 - e^2)^{3/2} \sqrt{\left[1 - 2 \frac{m}{l}(1 + e)\right] \left[1 - 2 \frac{m}{l}(1 - e)\right]} \int d\psi \left\{ \left(\frac{1 - e^2}{1 - e \cos \psi} \right)^{-2} \left(1 - 2 \frac{m}{l} \left(\frac{1 - e^2}{1 - e \cos \psi} \right) \right)^{-1} \left(1 - 2 \frac{m}{l} \left(2 + \frac{1 - e^2}{1 - e \cos \psi} \right) \right)^{-1/2} \frac{\sqrt{1 - e^2}}{1 - e \cos \psi} \right\}$$

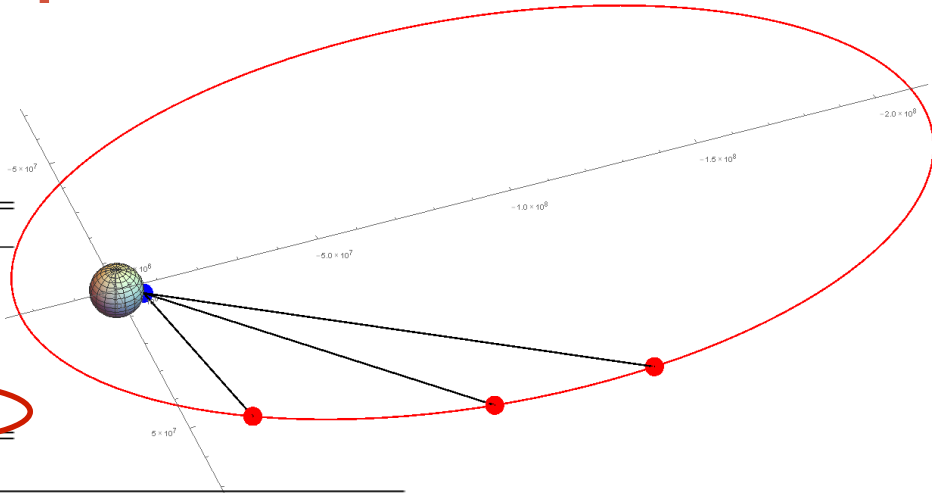
$$Y = 2\pi l \sqrt{\frac{l}{2m}} (1 - e^2)^{-3/2}$$

$$\phi = \int d\psi \left\{ \left[1 - 2 \frac{m}{l} \left(3 + \frac{e(e - \cos \psi)}{e \cos \psi - 1} \right) \right]^{-1/2} \frac{\sqrt{1 - e^2}}{1 - e \cos \psi} \right\}$$

Solving the inverse problem

- Elliptic orbits – 3 satellites

Order	t_P [s]	Error
Actual	1.0	
0	0.999981147292624478443271196316	1.8852×10^{-4}
1	1.000000000311447240513143137083	3.1144×10^{-10}
2	0.999999999999991902940522415661	8.0971×10^{-15}
3	1.0000000000000000240528310966	2.4052×10^{-19}



Order	r_P [10^6 m]	Error
Actual	6.3	
0	6.301820859016794837474852781689	2.8902×10^{-3}
1	6.299998869872663870009048889219	1.7938×10^{-7}
2	6.29999999725109007327543474487	4.3633×10^{-11}
3	6.30000000000019882419173483962	3.1559×10^{-15}

Order	ϕ_P [rad]	Error
Actual	2.617993877991494411361372840474	
0	2.617433575705034511049060737423	2.1401×10^{-3}
1	2.617993718180396219341713036414	6.1043×10^{-8}
2	2.617993878016665537549958414726	9.6147×10^{-12}
3	2.617993877991538123347549284001	1.6714×10^{-14}

Still to come

- Time transfer method applied to inclined eccentric orbits in solving the inverse problem
- Use Synge world function to produce null coordinates for eccentric orbits
- Determine whether Synge world function is applicable to the inverse problem
- Develop ABC system with these semi-analytic solutions
- Determine whether applicable to real world systems, e.g. the metric used in Andreja Gomboc's talk
- Compare with Andreja and Uros's numerical method
- Read off gravitational parameters
- Applications in relativistic astrophysics???